

SUPERCONDUCTIVITY AND VACANCY STRUCTURES OF THE PRESSURE-INDUCED
NaCl-TYPE PHASES OF THE In-Te SYSTEM

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(Received 23 December 1963)

We have found that a solid-solution range of Te exists in the pressure-induced NaCl-type InTe. The compounds are actually deficient in In; the formulas may be written In_{1-x}Te , with $0 \leq x \leq 0.18$. As far as we know, this is the first time structures with a substantial number of vacancies have been produced at high pressures. Superconductivity exists in the whole range of composition, the transition temperatures T_c decreasing with In/Te ratio, from 3.45–3.20°K for InTe to 1.06–1.02°K for $\text{In}_{0.82}\text{Te}$; the lattice constant decreases from 6.16 to 6.052 Å (see Table I and Fig. 1). An explanation

of the metallic behavior of these compounds has been deduced and appears also to account for the behavior of T_c with In/Te ratio.

Normal InTe undergoes transformation at 30 kbar to the NaCl-type phase¹ which is retained metastably at atmospheric pressure and room temperature. This phase was first reported to be superconducting at 2.18°K by Bömmel and others² but later Banus and others¹ reported a significantly higher transition temperature 3.7–3.5°K. Our result for InTe is in better agreement with the higher value.

We became interested in the NaCl-type phase when it occurred in some high pressure-high temperature experiments on In_2Te_3 and had a lattice constant of 6.052 Å, considerably smaller than that of stoichiometric InTe. Consequently several compositions in the range 33.3 to 55.0 at. % Te were prepared by melting together appropriate amounts of 99.99+ In and Te in sealed evacuated fused silica tubes which were continually agitated to ensure intimate mixing. X-ray powder photographs (CuK radiation) gave the expected patterns for normal In_2Te_3 and InTe. For compositions with In/Te ≤ 1 , mixtures of InTe and In_2Te_3 resulted as would be predicted

CONTENT ANALYSIS

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on	

Table I. Superconducting Transition Temperatures T_c and Lattice Constants a for In_{1-x}Te Compounds with NaCl-like Structure.

$1-x$	T_c ($^{\circ}\text{K}$)	a (\AA)
1.00	3.45 - 3.2	6.16 ± 0.01
0.95	2.7 - 2.5	6.14
0.91	2.04 - 1.87	6.110 ± 0.003
0.87	1.55 - 1.40	6.081
0.83	1.15 - 1.09	6.055
0.82	1.06 - 1.02	6.052

from the reported³ phase diagram. Pressure was generated in a piston cylinder device. The furnace assembly was the same as that described earlier;⁴ the sample container was a tantalum cup with a boron nitride cover. Pressure of about 30 kbar was applied first and the temperature raised to about 800 $^{\circ}\text{C}$. After a few minutes, the temperature was reduced to ambient in about 5 min and the pressure released. In experiments below the melting temperature, the specimens were kept under pressure at 550 $^{\circ}\text{C}$ for about one hour.

After being subjected to compression and heating, all specimens yielded metastably retainable NaCl-type phases. In the composition range InTe to $\text{InTe}_{1.22}$, single phases were obtained; In_2Te , however, transformed to a mixture of In and *stoichiometric* InTe , indicating (contrary to ref 1) that this phase does not dissolve excess In . The same results were obtained whether or not the specimens were melted.

The density of the NaCl-like $\text{In}_{0.82}\text{Te}$, determined pycnometrically, was 6.72 g/cc, which implies the formula $(\text{In}_{0.83}\text{Te}_{0.01}\square_{0.16})\text{Te}$; that is, each unit cell contains an average of 3.32 In , 0.04 Te and 0.64 vacancies in In sites. X-ray diffraction data give a clear corroboration of the vacancy structures. Because the difference in scattering by In and Te atoms is small, the $\{111\}$ line of InTe is very faint, but, as in the case of AgSbTe_2 ⁵ (and contrary to refs 1 and 2), can be seen. As In/Te decreases, the intensities of the odd-index reflections increase relative to the even-index ones. For $\text{In}_{0.82}\text{Te}$, the $\{111\}$, $\{311\}$, and $\{331\}$ reflections are readily seen. Weissenberg and Buerger precession camera photographs of a single crystal of $\text{In}_{0.82}\text{Te}$ corroborate the (statistical⁵) cubicity of this phase.

Tests for superconductivity on powdered specimens were made by the alternating current method of Schawlow and Devlin.⁶

We propose that the crystal structure of the high-pressure form of InTe is a disordered one (like that of AgSbTe_2 ⁵) containing equal numbers of In^+ and In^{3+} ions, and that its metallic behavior results from the ease with which electron transfer can occur from In^+ to In^{3+} ions throughout the crystal. Normal InTe (ref 7) has the TlSe (ref 8) structure and is a semiconductor. The TlSe crystal is ordered and contains equal numbers of Tl^+ and Tl^{3+} ions in 8 and 4 coordination respectively; consequently, it is expected that In^+ and In^{3+} ions are arranged similarly in normal InTe . The structure stabilizes the In valencies because it is energetically unfavorable for an In in 8 coordination to be trivalent or one in a tetrahedral site to be monovalent. The compound AgSbTe_2 ⁵ is isostructural with the pressure-induced form of InTe ; Ag^+ and Sb^{3+} ions have 6 coordination and their average valence is that of Te . In AgSbTe_2 transfer of electrons from Ag to Sb is also energetically unfavorable because the +1 valence state of Ag is highly stable; this compound is also a semiconductor.

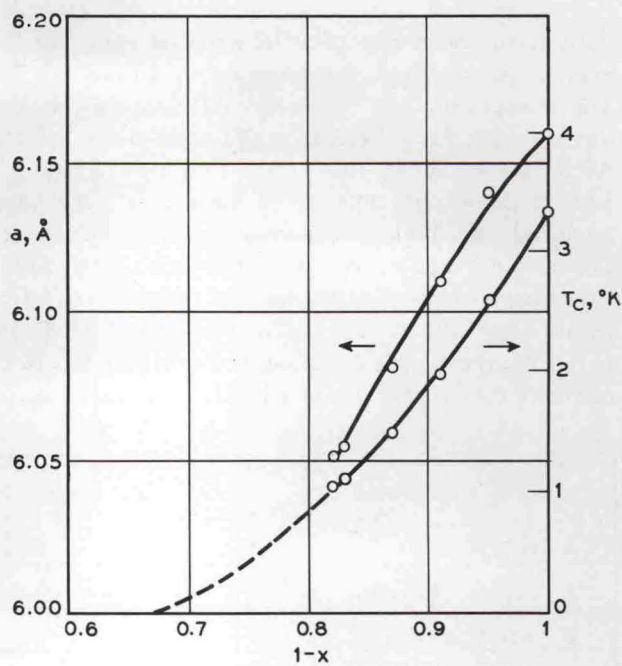


Fig. 1. Lattice constant a and superconducting transition temperature T_c vs $(1-x)$ for the system In_{1-x}Te .

As the ratio of In/Te decreases, the percentage of In^{3+} ions in the crystal increases. When In/Te = 0.82, the average valence of In is 2.44, implying that 72% of the In atoms are trivalent. It would appear that as the percentage of In^{3+} ions increases (with decreasing In/Te), the superconducting transition temperatures should decrease. The curve of T_c vs $(1-x)$ in Fig. 1, indicates that T_c does tend to approach 0 as $(1-x)$ approaches 0.67 at which point all In atoms would be trivalent.

We wish to thank T. H. Geballe for suggesting that the superconductivity tests be made on powdered specimens. We wish also to thank R. G. Maines for technical assistance.

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